Particle size dependence of magnetization and non-centrosymmetry in nanoscale BiFeO₃

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The saturation magnetization (M_S), antiferromagnetic transition point (T_N), and the off-center displacements of Bi and Fe ions have been measured as a function of particle size in nanoscale BiFeO₃. T_N decreases down to ~550 K for particles of size ~5 nm from ~653 K in bulk while M_S rises by more than an order of magnitude. Analysis of crystallographic structure from Rietveld refinement of x-ray diffraction patterns shows significant rise in off-center displacements of Bi (δ_{B_i}) and Fe (δ_{Fe}) ions within a unit cell with the decrease in particle size. The net unit-cell polarization P_S too, is found to be larger in nanoscale regime.

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I. INTRODUCTION

Extensive research work has been done on BiFeO₃ over the decades because of its magnetoelectric multiferroicity at room temperature.¹ The linear magnetoelectric coupling originates from striction mediated interaction between polar and magnetic domains.² The bulk $BiFeO_3$ is G-type antiferromagnetic with a spiral of long period (62) nm). Therefore, the magnetostriction is weak which, in turn, gives rise to weak coupling. The electric polarization is also small ($\sim 5-6 \ \mu C/cm^2$) in bulk.³ It improves $(\sim 55 \ \mu C/cm^2)$ in thin films due to the strain between substrate and film.⁴ Large polarization and magnetization will render BiFeO₃ quite useful as a single phase multiferroic compound. In this paper, we show that it is possible to achieve both the aspects - large polarization and magnetization - in nanoscale BiFeO₃ (\sim 5-50 nm). We found that intrinsic ferromagnetism develops in nanoscale $BiFeO_3$ with rise in saturation magnetization M_S , decrease in antiferromagnetic transition point T_N , and increase in coercive field H_C . We also found that the unit-cell off-center displacements of Bi and Fe ions which render BiFeO₃ ferroelectric - increases monotonically with the decrease in particle size. The net unit-cell polarization P_S too, is found be higher in nano-sized particles.

II. EXPERIMENTS

The experiments were carried out on both bulk and nano-sized particles of BiFeO₃. The nanoparticles of BiFeO₃ have been synthesized by sonochemical process where coprecipitation takes place from mixed aqueous solution of metal nitrates within a suitable medium in presence of ultrasonic vibration. The precipitate was washed in alcohol and dried. Finally, the product was calcined at 350° - 450° C for 2-6h. The sonochemical process was shown to yield nanosized particles of various oxide compounds in the past. We have separated out finer particles by a centrifuge running at different speeds - 10000, 12000, and 15000 rpm. The variation of heat-treatment temperature and time also yields particles of different sizes. In order to synthesize the bulk sample, the calcined powder was compacted in the form of pellets and sintered at 830° C for 5-10h. The average grain size in the sintered pellet was $\geq 0.1 \mu$ m.

The samples have been characterized by studying the x-ray diffraction (XRD) patterns at room temperature and from the transmission electron microscopy (TEM) and high resolution TEM. The electrical resistivity and magnetic measurements were carried out across a temperature range 300-900 K. The resistivity of the particles was measured by depositing a coating of the particles on an alumina substrate with gold electrodes printed on it. The XRD patterns were refined by Fullprof (ver 2.3) to determine the structural parameters such as space group, lattice parameters, crystallite size, strain, atom positions, bond lengths and angles etc. The off-center displacements of Bi³⁺ (δ_{Bi}) and Fe³⁺ (δ_{Fe}) ions have also been calculated as a function of particle size.

III. RESULTS AND DISCUSSION

The average size of the nanoscale BiFeO₃ particles varies from ~ 5 to ~ 50 nm. In Fig. 1, we show a representative TEM photograph. The average size was estimated by image analyzer software Image-J. It is compared with the crystallite size estimated from the Rietveld refinement of x-ray diffraction patterns. Both the results corroborate each other.

The magnetic measurements were carried out in zerofield cooled mode across 300-900 K under an applied field ~100 Oe. In Fig. 2a, we show the magnetization versus temperature plot for a few representative samples. Quite evident is the monotonic increase in magnetization with the decrease in particle size. This observation corroborates earlier observations by us as well as by other authors.^{5,6} Interestingly, the antiferromagnetic transition point T_N drops significantly with the decrease in particle size: from ~653 K in bulk to ~550 K in particles of average size ~5 nm. Dc electrical resistance has

TABLE I. List of lattice parameters, atom positions, bond length, angle, and microstrain at room temperature.

Samples	Lattice parameters (\mathring{A})	Atoms	Positions	х	у	Z	Bonds	length	Bonds	angle	microstrain	R_{wp}
BiFeO ₃ -bulk	a = 5.578 c = 13.868	Bi Fe O	6a 6a 18b	$0.0 \\ 0.0 \\ 0.4346$	0.0 0.0 0.0121	0.0 0.2198 -0.0468	Bi-O Fe-O Fe-O	2.309Å 1.949 Å 2.118 Å	Fe-O-Fe O-Bi-O	154.05^{o} 73.88^{o}	0.015	22.6
BiFeO ₃ - 25 nm	a = 5.573 c = 13.849	Bi Fe O	6a 6a 18b	$0.0 \\ 0.0 \\ 0.4715$	$0.0 \\ 0.0 \\ 0.0119$	0.0 0.22324 -0.0622	Bi-O Fe-O Fe-O	2.130 Å 1.804 Å 2.269 Å	Fe-O-Fe O-Bi-O	152.76° 78.93°	0.029	6.68
BiFeO ₃ - 22 nm	a = 5.579 c = 13.870	Bi Fe O	6a 6a 18b	$\begin{array}{c} 0.0 \\ 0.0 \\ 0.4455 \end{array}$	0.0 0.0 0.03123	0.0 0.2256 -0.05351	Bi-O Bi-O Fe-O Fe-O	2.179 Å 2.515 Å 1.955 Å 2.172 Å	Fe-O-Fe O-Bi-O	147.77^{o} 73.84^{o}	0.049	11.7
BiFeO ₃ - 19 nm	a = 5.624 c = 13.672	Bi Fe O	6a 6a 18b	$0.0 \\ 0.0 \\ 0.3712$	0.0 0.0 -0.0898	0.0 0.2285 0.00216	Bi-O Bi-O Fe-O Fe-O	2.381 Å 2.586 Å 1.941 Å 2.267 Å	Fe-O-Fe O-Bi-O	140.91^{o} 85.17^{o}	0.105	15.8



FIG. 1. TEM photographs for (a) finer and (b) coarser $BiFeO_3$ particles; inset shows the high resolution TEM photograph of a single crystalline nanoparticle.

also been measured across 300-800 K. Around T_N , a distinct feature in the resistance versus temperature plot could be noticed (data not shown here). T_N was thus estimated both from the magnetic and electrical measurements. In Fig. 2b, the T_N versus particle size plot is shown. There is a certain discrepancy between T_N estimated from magnetic and electrical measurements, especially, in finer particle regime. This could result from progressive broadening of the transition zone. We have carried out Curie-Weiss fitting of the magnetization data in the paramagnetic regime (Fig. 2a inset). With the decrease in particle size, the Weiss constant θ appears to be switching from negative (signaling antiferromagnetic order) to positive (marking ferromagnetism) regime. The degree of frustration $f = \theta/T_N$ increases from ~1.0 to ~ 1.4 with the decrease in particle size. This could be because of enhanced geometry-driven frustration in spin

ordering within near neighbors in finer particles. This has been observed in nanosized ferromagnetic particles by others as well.⁷ The Curie constant C yields the magnetization to be varying within ~0.108-0.452 μ_B /Fe for particle size varying within 5-50 nm. For the bulk system, the magnetization is ~0.02 μ_B /Fe. The room temperature saturation magnetization M_S and coercivity H_C are plotted as a function of particle size in Fig. 2b. The intrinsic ferromagnetism in nanoscale BiFeO₃ results from⁵ (i) incomplete spiral of magnetic order in particles of size less than ~62 nm, and (ii) enhanced strain of the nanocrystals.

The XRD patterns have been refined by Fullprof (ver 2.3) to find out the dependence of non-centrosymmetry of the crystallographic structure on particle size. In $BiFeO_3$, the ferroelectric polarization develops primarily because of covalency in Bi-O bonds resulting from Bi^{3+} 6s² lone pair and consequent off-center displacement of Bi^{3+} ions with respect to the near-neighbor oxygen network. Interestingly, the Fe^{3+} ions too occupy noncentrosymmetric positions in spite of Jahn-Teller distortion within $Fe^{3+}O_6$ octahedra. In Fig. 3a, we show the experimental and refined x-ray diffraction patterns for a representative case (~ 22 nm particle). Following points could be noted in the evolution of the patterns: (i) space group changes from $R\bar{3}c$ in bulk to lower symmetric R3c in the nanoscale samples; (ii) the strain increases with the decrease in particle size; (iii) the lattice parameters and volume change systematically with particle size. In Table-I, we provide the list of structural parameters as a function of particle size: (i) lattice parameters, (ii) atom positions, (iii) bond lengths and angles within a unit cell, and (iv) lattice strain. The weighted reliability factor \mathbf{R}_{wp} which defines the goodness of fit between experi-



FIG. 2. (color online) (a) Magnetic moment versus temperature plots for a few representative samples of different particle size; the magnetic transition near T_N could be clearly observed in bulk and coarser particles; inset shows the Curie-Weiss plot in the paramagnetic regime; (b) variation in the magnetic transition point T_N as estimated from magnetic (open circle) and electrical (up-triangle) measurements; the variation of M_S (solid circle) and H_C (diamond) with particle size are also shown; inset shows the room temperature hysteresis loops for bulk (open circle) and nanoparticles of size 25 nm (up-triangle) and 40 nm (down-triangle).

mental and refined XRD patterns is also shown.

We estimated the off-center displacements of Fe³⁺ and Bi³⁺ ions within a unit cell of BiFeO₃ from the atom position data. In bulk sintered pellet, the δ_{Bi} and δ_{Fe} are ~1.235 Å and ~0.229 Å, respectively. Remarkably, both δ_{Bi} and δ_{Fe} exhibit monotonic rise as the particle size is reduced (Fig. 3b).⁸ Interestingly, in contrast to the results in Ref. 8, we found that while δ_{Fe} results in polarization along $[111]_{rh} \parallel [001]_{hex}$ axis, δ_{Bi} is oriented in a different direction which gives polarization along all three directions. Using the δ_{Fe} and the component of δ_{Bi} along $[111]_{rh} \parallel [001]_{hex}$, we estimated the polarization P_S in a unit cell following the procedure described in Ref. 9. P_S is plotted in Fig. 3b as a function particle size. It shows that the resultant P_S of a unit cell increases substantially in nanoscale regime around ~28



FIG. 3. (color online) (a) The experimental and Rietveld refined x-ray diffraction patterns for a representative case (~ 22 nm particle); inset shows the corresponding crystallographic structure (rhombohedral with R3c space group); big, medium, and small spheres represent Bi, Fe, and O ions, respectively; (b) variation of off-center displacements of Bi and Fe ions as a function of particle size in the nanoscale regime; variation of net unit cell polarization P_S along the $[111]_{rh} \parallel [001]_{hex}$ with particle size is also shown.

nm. It decreases with further decrease in particle size and eventually stabilizes at a value which is much higher than what has been observed in bulk sample³ or even in thin films.⁴ This is the central result of this paper. Of course, the large unit cell polarization may not yield large ferroelectric polarization for the entire nanoparticle. This is because of the depolarization originated from surface effects (dead layer) of a nanoparticle¹⁰. It is necessary to design an appropriate electrode-nanoparticle architecture to cancel out the depolarizing field and thus observe large polarization even across an entire nano-sized particle.

IV. CONCLUSION

We observe significant improvement in unit cell polarization in nanoscale $BiFeO_3$. Both the off-center displacements of Bi^{3+} and Fe^{3+} ions increase monotonically with the decrease in particle size. The nanosized particles also exhibit intrinsic ferromagnetic order. Large magnetostriction resulting from ferromagnetism, possibly, gives rise to striction mediated enhanced ferroelectric polarization in a unit cell. Large ferroelectric polarization and ferromagnetism in nanosized particles will render nanoscale $BiFeO_3$ even more useful than the bulk

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system for magnetoelectric device applications.

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